

# 1 Mathematical Formulation

I prefer to use

$$\exp\{i(\omega t - \beta z)\}$$

for the waves in  $z$  direction. Since all equations will be in the frequency domain, in the following text  $\exp(i\omega t)$  will be omitted, and the symbol  $t$  will be re-used for the function parameter, instead of time. In addition,  $\epsilon = \epsilon_0$  and  $\mu = \mu_0$  for convenience.

From the two Maxwell equations inside the waveguide

$$\nabla \times \mathbf{H} = i\omega \epsilon \mathbf{E} \quad (1a)$$

$$\nabla \times \mathbf{E} = -i\omega \mu \mathbf{H} \quad (1b)$$

One write the Helmholtz equation using  $\mathbf{U} \in \{\mathbf{E}, \mathbf{H}\}$

$$\nabla^2 \mathbf{U} + \omega^2 \mu \epsilon \mathbf{U} = 0 \quad (2)$$

Consider first only the ideal electric conductor ( $\sigma \rightarrow \infty$ ) for waveguide walls

$$\mathbf{n} \times \mathbf{E} = 0 \quad (3a)$$

$$\mathbf{n} \cdot \mathbf{B} = 0 \quad (3b)$$

Let  $z = 0$  be the reference plane.  $(x', y')$  is a point in this plane, the point has the azimuthal angle

$$\xi := \arg(x', y') \quad (4)$$

The “trajectory” of a point in the helical waveguide is a helix at radius

$$r := \sqrt{x'^2 + y'^2} \quad (5)$$

A given tuple  $(x', y')$  or  $(r, \xi)$  labels one of the helical trajectories in Cartesian or cylindrical coordinates, respectively.

A point at  $(x', y', 0)$  will follow the helical trajectory to an arbitrary  $z$ . This process is a counterclockwise affine rotation by angle  $\iota z$ , where the constant  $\iota$  is the pitch of the helix

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \underbrace{\begin{pmatrix} \cos(\iota z) & -\sin(\iota z) & 0 \\ \sin(\iota z) & \cos(\iota z) & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{= M(z)} \begin{pmatrix} x' \\ y' \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ z \end{pmatrix} \quad (6a)$$

The rotation matrix is denoted as  $M$ . To obtain the label of helix, the inverse transformation can be performed

$$\begin{pmatrix} x' \\ y' \\ 0 \end{pmatrix} = \underbrace{\begin{pmatrix} \cos(\iota z) & \sin(\iota z) & 0 \\ -\sin(\iota z) & \cos(\iota z) & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{= M^{-1}(z) = M(-z)} \begin{pmatrix} x \\ y \\ z \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ -z \end{pmatrix} \quad (6b)$$

Hence,

$$\frac{\partial x'}{\partial x} = \cos(\iota z) \quad \text{and} \quad \frac{\partial x'}{\partial y} = \sin(\iota z) \quad (7a)$$

$$\frac{\partial y'}{\partial x} = -\sin(\iota z) \quad \text{and} \quad \frac{\partial y'}{\partial y} = \cos(\iota z) \quad (7b)$$

and the second derivatives of these are zero.

The key for a wave propagating (or whatever) in a rotating  $z$  direction is the Ansatz of variable pseudo-separation

$$U(x, y, z) = M(z) \mathbf{u}(x', y') \exp(-i \beta z) \quad (8)$$

Note that the transformation for  $z$  component is identity. For an easier interpretation, one can read

$$\mathbf{u}(x', y') = U(x', y', 0)$$

Since the transformation  $M(z)$  and the  $\exp(\dots)$  term is independent of  $x$  and  $y$

$$\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} = M(z) \left( \frac{\partial^2 \mathbf{u}}{\partial x^2} + \frac{\partial^2 \mathbf{u}}{\partial y^2} \right) \exp(-i \beta z)$$

The derivative of  $z$  is

$$\frac{\partial \mathbf{U}}{\partial z} = \left\{ \frac{\partial M}{\partial z} \mathbf{u} + M \frac{\partial \mathbf{u}}{\partial z} - \mathbf{i} \beta M \mathbf{u} \right\} \exp(-\mathbf{i} \beta z)$$

and the second-order derivative is

$$\begin{aligned} \frac{\partial^2 \mathbf{U}}{\partial z^2} = & \left\{ \frac{\partial^2 M}{\partial z^2} \mathbf{u} + 2 \frac{\partial M}{\partial z} \left( \frac{\partial \mathbf{u}}{\partial z} - \mathbf{i} \beta \mathbf{u} \right) \right. \\ & \left. + M \left( \frac{\partial^2 \mathbf{u}}{\partial z^2} - 2 \mathbf{i} \beta \frac{\partial \mathbf{u}}{\partial z} - \beta^2 \mathbf{u} \right) \right\} \exp(-\mathbf{i} \beta z) \end{aligned}$$

$$\begin{aligned} \nabla^2 \mathbf{U} = \frac{\partial^2 \mathbf{U}}{\partial x^2} + \frac{\partial^2 \mathbf{U}}{\partial y^2} + \frac{\partial^2 \mathbf{U}}{\partial z^2} = & \left\{ \frac{\partial^2 M}{\partial z^2} \mathbf{u} + 2 \frac{\partial M}{\partial z} \left( \frac{\partial \mathbf{u}}{\partial z} - \mathbf{i} \beta \mathbf{u} \right) \right. \\ & \left. + M \left( \nabla^2 \mathbf{u} - 2 \mathbf{i} \beta \frac{\partial \mathbf{u}}{\partial z} - \beta^2 \mathbf{u} \right) \right\} \exp(-\mathbf{i} \beta z) \end{aligned} \quad (9)$$

Treat the Helmholtz equation in eq. (2) as

$$M M^{-1} \nabla^2 \mathbf{U} + \omega^2 \mu \epsilon \mathbf{U} = 0 \quad (10)$$

then the transformation  $M$  ahead of each term and the  $\exp(-\mathbf{i} \beta z)$  part can be cancelled

$$\begin{aligned} M^{-1} \frac{\partial^2 M}{\partial z^2} \mathbf{u} + 2 M^{-1} \frac{\partial M}{\partial z} \left( \frac{\partial \mathbf{u}}{\partial z} - \mathbf{i} \beta \mathbf{u} \right) \\ + \nabla^2 \mathbf{u} - 2 \mathbf{i} \beta \frac{\partial \mathbf{u}}{\partial z} + (\omega^2 \mu \epsilon - \beta^2) \mathbf{u} = 0 \end{aligned} \quad (11)$$

It can be found that

$$M^{-1} \frac{\partial^2 M}{\partial z^2} = \begin{pmatrix} -\iota^2 & 0 & 0 \\ 0 & -\iota^2 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (12a)$$

and

$$M^{-1} \frac{\partial M}{\partial z} = \begin{pmatrix} 0 & -\iota & 0 \\ \iota & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (12b)$$

Note that  $\iota = 0$  should in agreement with the usual variable separation in a non-rotating coordinates. (Hint: the terms with  $M$  vanish. Besides,  $x'$  and  $y'$  are in this case independent of  $z$ , thus  $\partial \mathbf{u} / \partial z = 0$ .)

$$\begin{aligned}
\frac{\partial^2 \mathbf{u}}{\partial x^2} &= \cos(\iota z) \frac{\partial}{x} \frac{\partial \mathbf{u}}{\partial x'} - \sin(\iota z) \frac{\partial}{x} \frac{\partial \mathbf{u}}{\partial y'} \\
&= \cos(\iota z) \left( \cos(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x'^2} - \sin(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x' \partial y'} \right) \\
&\quad - \sin(\iota z) \left( \cos(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x' \partial y'} - \sin(\iota z) \frac{\partial^2 \mathbf{u}}{\partial y'^2} \right) \\
&= \cos^2(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x'^2} + \sin^2(\iota z) \frac{\partial^2 \mathbf{u}}{\partial y'^2} - 2 \sin(\iota z) \cos(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x' \partial y'}
\end{aligned}$$

For the twice derivative of  $y$ , since there is no transformation of triangle functions in the last steps, it is just to replace wherever  $\cos$  by  $-\sin$  while  $\sin$  by  $\cos$ :

$$\frac{\partial^2 \mathbf{u}}{\partial y^2} = \sin^2(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x'^2} + \cos^2(\iota z) \frac{\partial^2 \mathbf{u}}{\partial y'^2} + 2 \sin(\iota z) \cos(\iota z) \frac{\partial^2 \mathbf{u}}{\partial x' \partial y'}$$

Therefore,

$$\nabla_{\perp}^2 \mathbf{u} = \frac{\partial^2 \mathbf{u}}{\partial x^2} + \frac{\partial^2 \mathbf{u}}{\partial y^2} = \frac{\partial^2 \mathbf{u}}{\partial x'^2} + \frac{\partial^2 \mathbf{u}}{\partial y'^2} = \nabla_{\perp}'^2 \mathbf{u}$$

The calculation of  $z$  derivatives is tricky. First, from eq. (6b)

$$\frac{\partial}{\partial z} \begin{pmatrix} x' \\ y' \end{pmatrix} = \iota \begin{pmatrix} -\sin(\iota z) & \cos(\iota z) \\ -\cos(\iota z) & -\sin(\iota z) \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

substituting  $(x, y)$  by eq. (6a)

$$\frac{\partial}{\partial z} \begin{pmatrix} x' \\ y' \end{pmatrix} = \iota \begin{pmatrix} -\sin(\iota z) & \cos(\iota z) \\ -\cos(\iota z) & -\sin(\iota z) \end{pmatrix} \begin{pmatrix} \cos(\iota z) & -\sin(\iota z) \\ \sin(\iota z) & \cos(\iota z) \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix}$$

that is

$$\frac{\partial}{\partial z} \begin{pmatrix} x' \\ y' \end{pmatrix} = \iota \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \iota y' \\ -\iota x' \end{pmatrix} \quad (13)$$

with this shortcut,  $\partial \mathbf{u} / \partial z$  can be write as

$$\frac{\partial \mathbf{u}(x', y')}{\partial z} = \frac{\partial \mathbf{u}}{\partial x'} \frac{\partial x'}{\partial z} + \frac{\partial \mathbf{u}}{\partial y'} \frac{\partial y'}{\partial z} = \iota \left( y' \frac{\partial \mathbf{u}}{\partial x'} - x' \frac{\partial \mathbf{u}}{\partial y'} \right) \quad (14a)$$

The second-order derivative is

$$\frac{\partial^2 \mathbf{u}(x', y')}{\partial z^2} = \iota^2 \left[ y'^2 \frac{\partial^2 \mathbf{u}}{\partial x'^2} + x'^2 \frac{\partial^2 \mathbf{u}}{\partial y'^2} - \left( x' \frac{\partial \mathbf{u}}{\partial x'} + y' \frac{\partial \mathbf{u}}{\partial y'} + 2 x' y' \frac{\partial^2 \mathbf{u}}{\partial x' \partial y'} \right) \right] \quad (14b)$$

Equation (10) can be written in  $(x', y')$  coordinates

$$\begin{aligned} & -\iota^2 \begin{pmatrix} u_x \\ u_y \\ 0 \end{pmatrix} + 2 \iota y' \frac{\partial}{\partial x'} \begin{pmatrix} -u_y \\ u_x \\ 0 \end{pmatrix} + 2 \iota x' \frac{\partial}{\partial y'} \begin{pmatrix} u_y \\ -u_x \\ 0 \end{pmatrix} + 2 \text{i} \beta \begin{pmatrix} u_y \\ -u_x \\ 0 \end{pmatrix} \\ & + (1 + \iota^2 y'^2) \frac{\partial^2}{\partial x'^2} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} + (1 + \iota^2 x'^2) \frac{\partial^2}{\partial y'^2} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} \\ & - \iota^2 x' \frac{\partial}{\partial x'} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} - \iota^2 y' \frac{\partial}{\partial y'} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} - 2 \iota^2 x' y' \frac{\partial^2}{\partial x' \partial y'} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} \\ & - 2 \text{i} \beta \iota y' \frac{\partial}{\partial x'} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} + 2 \text{i} \beta \iota x' \frac{\partial}{\partial y'} \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} + (\omega^2 \mu \epsilon - \beta^2) \begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = 0 \end{aligned}$$